

# Chemical Crystallography and Structural Chemistry

VO 270287

Lecture N° 6 — 19<sup>th</sup> May 2022

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## Course Details

3 <sup>rd</sup>	March	Lecture N <sup>o</sup> 1	10 <sup>th</sup>	March	Lecture N <sup>o</sup> 2
17 <sup>th</sup>	March	<b>no lecture</b>	24 <sup>th</sup>	March	Exercise N <sup>o</sup> 1
31 <sup>st</sup>	March	Lecture N <sup>o</sup> 3			
7 <sup>th</sup>	April	Discussion Ex. 1	14 <sup>th</sup>	April	<b>Easter break</b>
21 <sup>st</sup>	April	<b>Easter break</b>	28 <sup>th</sup>	April	Lecture N <sup>o</sup> 4
5 <sup>th</sup>	May	Power cut	12 <sup>th</sup>	May	<b>no lecture</b>
19 <sup>th</sup>	May	Lecture N <sup>o</sup> 6	26 <sup>th</sup>	May	Exercise N <sup>o</sup> 3
2 <sup>nd</sup>	June	Lecture N <sup>o</sup> 7	8 <sup>th</sup>	June	Lecture N <sup>o</sup> 8
16 <sup>th</sup>	June	Exercise N <sup>o</sup> 4	23 <sup>th</sup>	June	Lecture N <sup>o</sup> 9
30 <sup>th</sup>	June	<b>no lecture</b>			

## Previous Lecture

- Symmetry and Diffraction data
- 11 Laue groups
- Friedel's law
- Some aspects on data collection
- Indexing

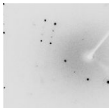
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# 1 Experimental procedure: from data collection to structure

## From Data Collection to Structure

**Data collection**



several GB



**Data integration**

0	0	-1	2.7	0.9
0	0	1	4.0	1.0
0	0	-2	1'257.0	35.5
0	0	-2	1'600.0	42.7

several files, 100's MB



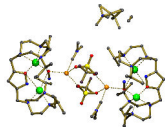
**Data Scaling**

0	0	-1	2.8	0.55
0	0	1	3.8	0.63
0	0	-2	1'432.0	95.7
0	0	-2	1'282.0	85.9

1 "hkl"-file, 50MB



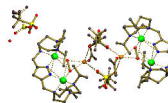
**Phasing**



Starting model



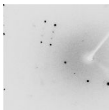
**Refinement**



Chemically  
sensible  
model

## Data Collection

Data collection



several GB

Data integration

0	0	-1	2.7	0.9
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several files, 100's MB

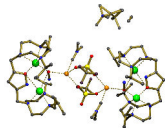
Data Scaling

0	0	-1	2.8	0.55
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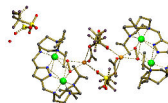
Phasing



Starting model



Refinement



Chemically  
sensible  
model

## How to collect good data

Data resolution

Data completeness

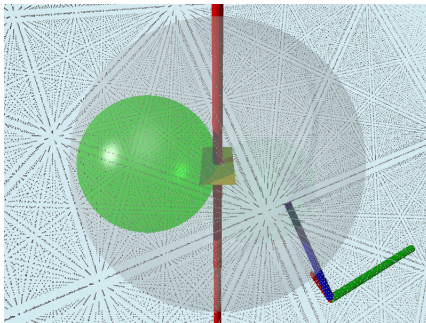
Data multiplicity

Data quality



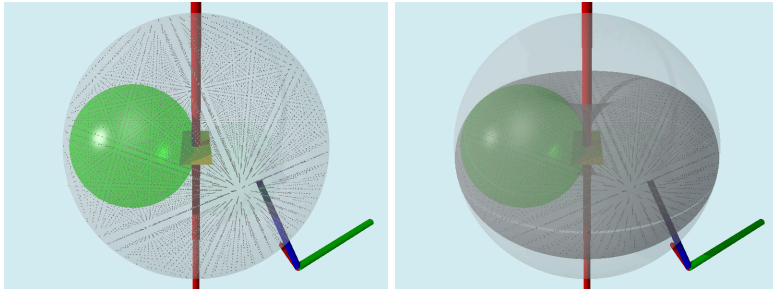
## Data resolution

- highest possible resolution:  $d_{\min} > \lambda/2$
- good crystals diffract better than this limit (instrumentation limited resolution)
- bad crystals may not reach this resolution (sample limited resolution)
- NB: low resolution are also important



## Data completeness

- Only reflections with resolution sphere can be collected (radius in reciprocal space:  $r = 2./\lambda$ )
- Laue equations permit to compute their location, and how to rotate the crystal into an appropriate position.
- Single rotation axis: torus shape of reflections.

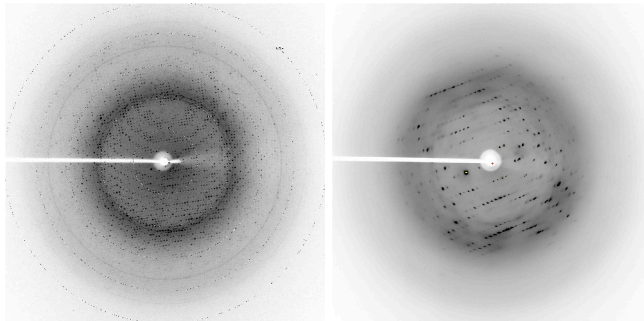


## Data multiplicity

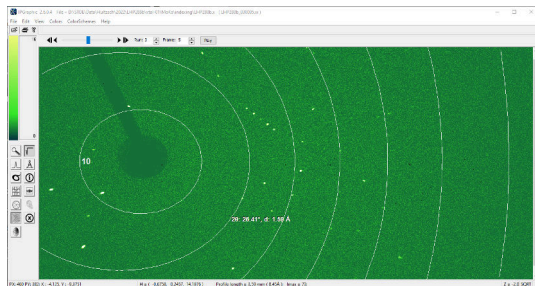
- Intensity ( $I(hkl)$ ) of the identical reflection ( $hkl$ ) can be measured multiple times
- Intensities  $I(hkl)$  and  $I'(R(hkl))$  symmetry equivalent reflections increase multiplicity
- Measuring on different positions on detector: reduction of systematic errors
- increasing multiplicity reduces uncertainty  $\sigma(I)$

## Data quality

- “Diffraction data acquisition is the final experimental stage of the crystal structure analysis.” (Z. Dauter, [1])
- Ensure proper data collection strategy (desired resolution, multiplicity; adjust detector distance to unit cell parameters, etc.)
- Crystal quality can be the limiting factor

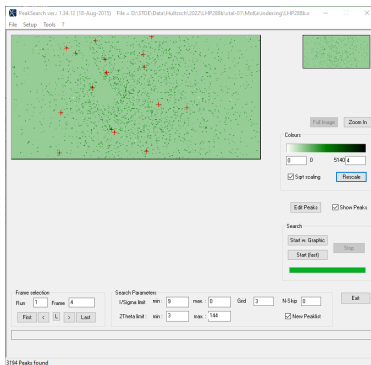


## Steps of data integration: initial diffraction images



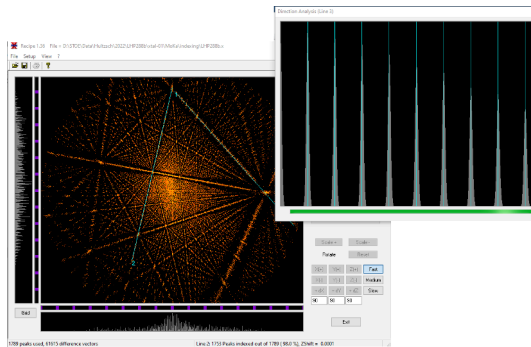
- typically four short runs:
- $2 \times 10^\circ - 15^\circ$ :  $0^\circ - 10^\circ$  and  $90^\circ - 100^\circ$
- two detector settings each (low to mid resolution and mid to high resolution)

## Steps of data integration: peak finding



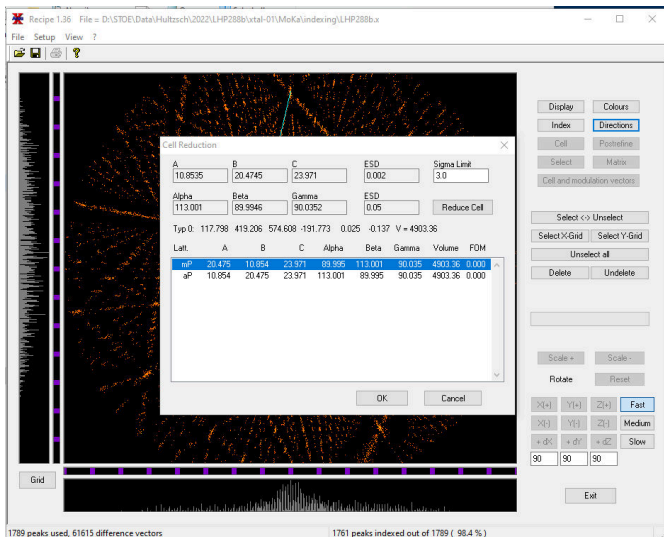
- simple (non-crystallographic) pattern recognition algorithm: peak above background
- main parameter: desired peak strength ( $I/\sigma$ )

## Steps of data integration: finding unit cell



- known instrument geometry: project peaks into reciprocal space
- find a lattice basis

## Steps of data integration: finding unit cell



Recipe 1.36 File = D:\STOE\Data\Hultzschi,2022\LHP288b\xtal-01\MoKa\indexing\LHP288b.x

File Setup View ?

Cell Reduction

A	B	C	ESD	Sigma Limit
10.8535	20.4745	23.971	0.002	3.0
Alpha	Beta	Gamma	ESD	
113.001	89.9946	90.0352	0.05	Reduce Cell

Typ 0: 117.798 419.206 574.608 -191.773 0.025 -0.137 V = 4903.36

Latt.	A	B	C	Alpha	Beta	Gamma	Volume	FOM
hP	20.475	10.854	23.971	89.995	113.001	90.035	4903.36	0.000
aP	10.854	20.475	23.971	113.001	89.995	90.035	4903.36	0.000

1789 peaks used, 61615 difference vectors

1761 peaks indexed out of 1789 ( 98.4 % )

Grid

Display Colours

Index **Directions**

Cell Postrefine

Select Matrix

Cell and modulation vectors

Select <> Unselect

Select X-Grid Select Y-Grid

Unselect all

Delete Undo

Scale + Scale -

Rotate Reset

X(-) Y(-) Z(-) Fast

X(+) Y(+) Z(+) Medium

+dX +dY +dZ Slow

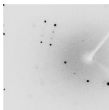
90 90 90

Exit



## 2 Data Integration

Data collection



several GB



Data integration

0	0	-1	2.7	0.9
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several files, 100's MB



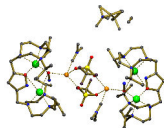
Data Scaling

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1 "hkl"-file, 50MB



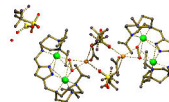
Phasing



Starting model



Refinement



Chemically  
sensible  
model

## Data Integration

Data integration comprises

1. Indexing: Determination of unit cell dimensions, orientation of the crystal, point group
2. Extraction of spot intensities from detector images.
3. Optimisation of experimental parameters

## Programs for data integration (incomplete)

**Saint** Licensed by Bruker AXS. Specific to Bruker programs. Good for data from twinned crystals. Derived from XDS.

**X-Area** Software suite distributed with Stoe X-ray diffractometers. Free to use. Very good interactive indexing routine, good for twinned data.

**XDS** Free for non-commercial users (<http://xds.mpimf-heidelberg.mpg.de>). Supports nearly all detector formats, very well documented. Very fast.

**DIALS** Free for non-commercial users. Very active development. (<https://dials.diamond.ac.uk/>)

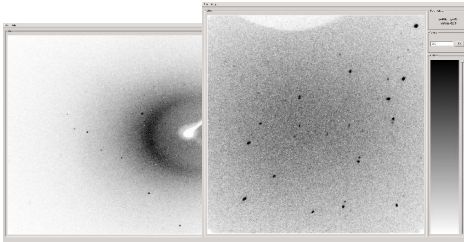
**EVAL Suite** Free for non-commercial users (<http://www.crystal.chem.uu.nl/distr/eval>). Can integrate e.g. incommensurate crystals

**Crysalis Pro** Licensed by Rigaku (<https://www.rigaku.com/products/smc/crysalis>)

**iMosflm** Free for non-commercial users, distributed with CCP4 (<http://www.ccp4.ac.uk>)

**HKL3000** Very good visualisation GUI for fine-tuning of parameters. very popular in the US (<https://hkl-xray.com/>)

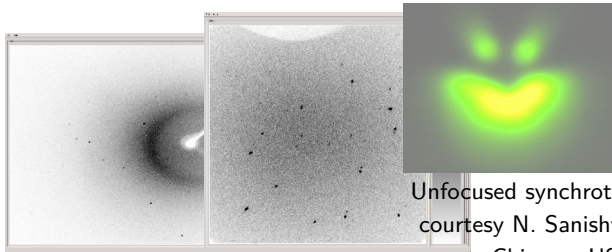
## Indexing



“Ewald sphere backwards:”

- find 200–1000 **strong** spots
- backtransform into reciprocal space (Laue equations)
- find a lattice and a suitable basis
- basis for reciprocal lattice corresponds to reduced unit cell constants

## Indexing



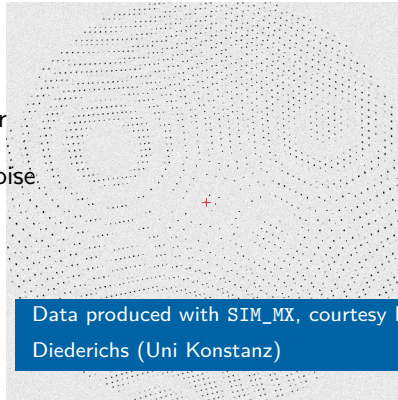
Unfocused synchrotron beam,  
courtesy N. Sanishvili, APS,  
Chicago, USA

Possible reasons for indexing problems:

- Incorrect parameters: detector distance, direction of rotation, wavelength, especially at synchrotrons
- Too few reflections
- Distorted spots (lattice defects, unfocused beam)
- Alien spots (ice, metal, contaminant)
- multiple lattices *twins*

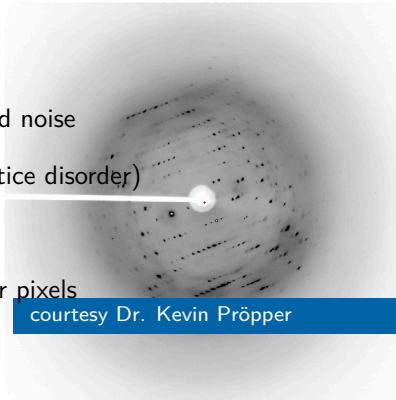
## Ideal diffraction image

- no lattice disorder
- no background noise
- perfect beam



## Realistic diffraction image

- strong background noise
- smeary spots (lattice disorder)
- spot overlap
- saturated detector pixels
- finite resolution



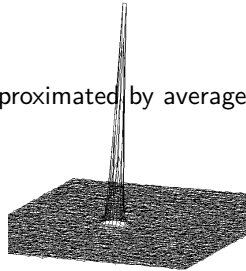


## Signal extraction

1. Calculate reflex positions
2. Determine local background
3. Differentiate background from signal (spot volume, shape)
4. Different approach for strong spots and weak spots

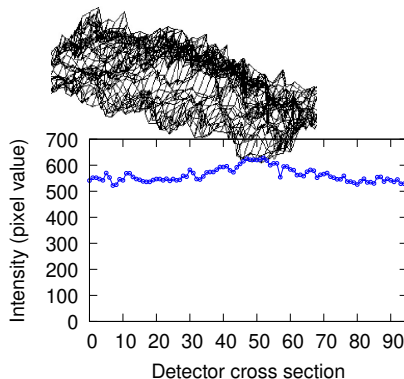
## Signal extraction: strong spots

- High intensity  $\Rightarrow$  Small error from noise
- Spot covers large detector area: noise approximated by average
- Good spot separation
- Good spot profile (shape)

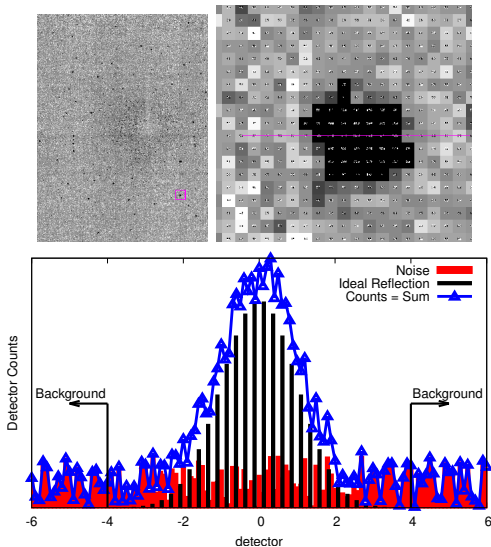


## Signal extraction: weak spots

- Low intensity:  $\Rightarrow$  high error from noise
- Spot covers small detector area: large effect from noise
- Background difficult to determine



## Cross section of a spot on the detector



## Reflection profile

Most integration programs create a set of reflection profiles from strong and reliable reflections, e.g. a 3D Gauss function.

The profiles depend on the region on the detector and on the crystal orientation.

Advantages:

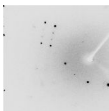
1. Measurement of weak reflections (fitting data for profile)
2. takes non-isotropic crystal shape into account
3. takes regions of varying detector sensitivity into account
4. produces a standard deviation of the reflection intensity:  $(h, k, l, I, \sigma_I)$

## Summary Data Integration

- Starts with indexing: crystal orientation, unit cell
- Look at all images per run
- Look only at calculated spot positions on detector
- Strong spots: sum pixel values, subtract background
- Strong spots: determine average reflection profile
- Weak spots: extract data based on profiles

### 3 Scaling

Data collection



several GB



Data integration

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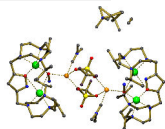
Data Scaling

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1 "hkl"-file, 50MB



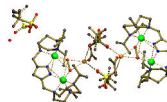
Phasing



Starting model



Refinement



Chemically  
sensible  
model

## Objective of the diffraction experiment

Structure elucidation of a chemical compounds

- chemical composition (e.g. purity after chromatography)
- Connectivity, distances between (non-) bonded atoms
- Configuration of stereochemical centres (R,S)

The values should be independent from the experimental setup.

Scaling makes the raw intensities (from data integration) independent from the experimental setup

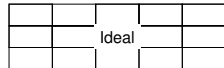


## Calculation of Intensities

Under consideration of the experimental, non-idealised setup, intensities are calculated as [2, Ch. Diffraction Intensities]

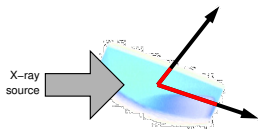
$$I_{exp}(hkl) = \frac{e^4}{m_e^2 c^4} \frac{\lambda^3 V_{crystal}}{V_{u.c.}^2} I_0 L P T E |F_{theor.}(hkl)|^2 \quad (1)$$

- $I_0$  incoming intensity (may vary with time)
- $L$  Lorentz factor describes trajectory through the Ewald sphere
- $P$  Polarisation correction;  $P = (1 + \cos^2 2\theta)/2$  for unpolarised source
- $T$  Absorption correction (esp. heavy elements)
- $E$  extinction correction (crystal defects, mosaicity)
- $F_{theor.}(hkl)$  structure factor (calculated from structure)

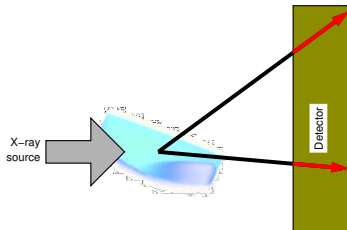


## Scaling = Idealisation and correction

In order to make data as independent from the experiment as possible, data are “standardised”. Some corrections are of numerical nature (polarisation), others are sample dependent (extinction: depends on elements in compounds). Two examples: absorption and angle of incidence



Absorption in the crystal depends on the path



Higher angle of incidence w.r.t. detector surface leads longer path through detector phosphor and thus to stronger signal

## Basis for Scaling: Symmetry and multiple measurements

- Some corrections depend on the instrument and can be calibrated (polarisation, angle of incidence, Lorentz factor)
- Some corrections (e.g. absorption, extinction) are (also) sample dependent

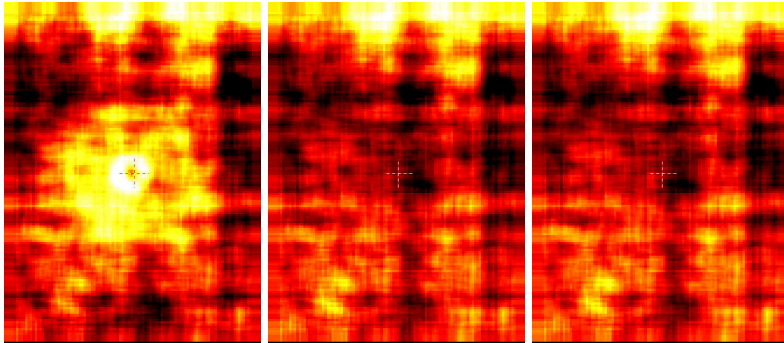
**Idea:** symmetry equivalent reflections, or multiply measured reflections, should have the same intensity

### Scaling means

1. Determination of the measured intensity  $I_{obs}$  of a set of equivalent reflections
2. Determination of their standard uncertainties
3. Result: idealised data set

## Examples for corrections

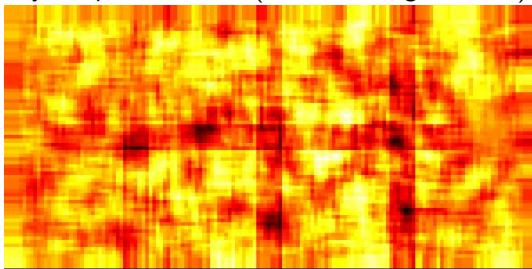
CMOS Detector (Bruker Photon II)



Detector “gain”: dark regions of the detector are more sensitive (factor between 0.97 and 1.3). Direct beam (left): indicates non-linear response at high intensity

## Examples for corrections

Hybrid pixel detector (DECTRIS Eiger2 R500)

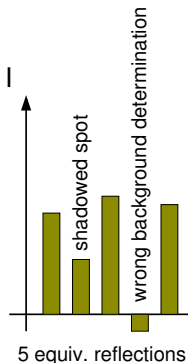


Hybrid pixel detectors are closer to ideal detectors for X-rays, with a smooth gain (factor between 0.96 and 1.02).

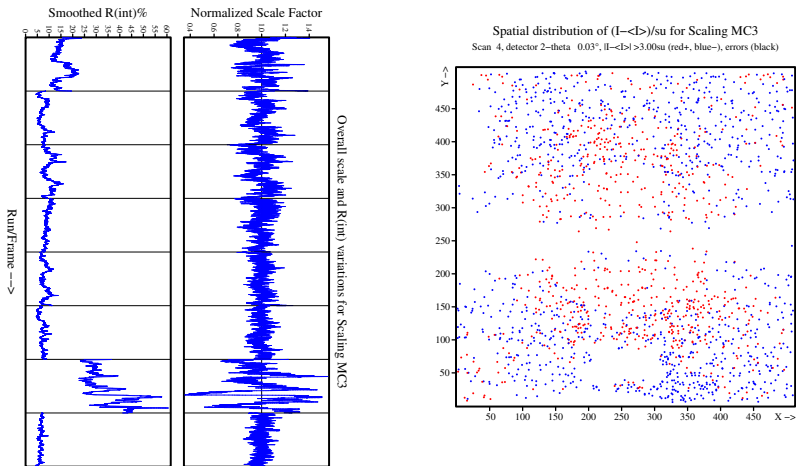
## Scaling details

Options to find the mean intensity  $\langle I \rangle$  of a group of equivalent reflections

1. Statistical average: exaggerates outliers
2. Scaling of each group of equivalents independently from other groups: neglects systematic errors.
3. Instead: **One** scale factor for several groups of equivalent reflections (e.g. 100 groups) which are close together on the detector



## Scaling plots (program SADABS)



## Idealised intensities

Before Scaling

After Scaling

$$I(hkl) = \frac{e^4}{m_e^2 c^4} \frac{\lambda^3 V_{\text{crystal}}}{V_{\text{u.c.}}^2} I_0 L P T E |F(hkl)|^2 \quad I(hkl) = c |F(hkl)|^2 \quad (2)$$

0	-1	5	1.379E+03	2.516E+02	0	-1	5	7.014E+01	1.208E+01
0	-1	-5	1.367E+03	2.726E+02	0	-1	-5	6.812E+01	1.274E+01
0	1	5	1.184E+03	2.610E+02	0	1	5	5.987E+01	1.231E+01
0	1	-5	1.347E+03	2.674E+02	0	1	-5	6.753E+01	1.258E+01
0	-1	6	1.090E+04	-1.229E+03				outlier removed	
0	-1	-6	4.677E+03	5.733E+02	0	-1	-6	2.365E+02	2.856E+01
0	1	6	4.286E+03	5.488E+02	0	1	6	2.145E+02	2.689E+01
0	1	-6	9.065E+03	-1.034E+03				outlier removed	
0	-1	7	0.204E+02	0.571E+01	0	-1	7	1.404E+02	2.271E+01



## Summary Scaling

- Scaling finds the mean intensity  $\langle I \rangle$  of symmetry equivalent reflections
- Scaling puts one or more data sets on a common scale
- Scaling removes outliers
- The result of scaling is a single  $(hkl)$ -file, which is used for all subsequent steps:
  1. Phasing, *i.e.* structure solution
  2. Model refinement

## 4 Intermezzo: The structure factor

## The Structure Factor $F(hkl)$

- Context between atoms and diffraction intensities
- Describing the electron density with the independent atom model (“IAM”)
- Formfactor and the “fudge factor” ADP

## The Structure Factor

The structure factor  $F(hkl)$  is related to the **electron density**  $\rho(x, y, z)$ , i.e. the distribution of electrons inside the unit cell:

$$F(hkl) = \int_{\text{unit cell}} \rho(x, y, z) e^{2\pi i(hx+ky+lz)} d^3x \quad (3)$$

This equation is the Fourier transformation of the electron density.

*Note:* The term “Fourier transformation” is important mainly because computers are very fast in calculating Fourier transformations.

## The Structure Factor

The *Independent Atom Model* (IAM, alias *isolated atom model*) is a powerful method to calculate the atomic structure factor  $F(hkl)$

$$F(hkl) = \sum_{\substack{\text{atoms } j \\ \text{in u.c.}}} f_j(\theta) e^{-8\pi^2 U_j(\theta, \lambda)} e^{2\pi i(hx_j + ky_j + lz_j)} \quad (4)$$

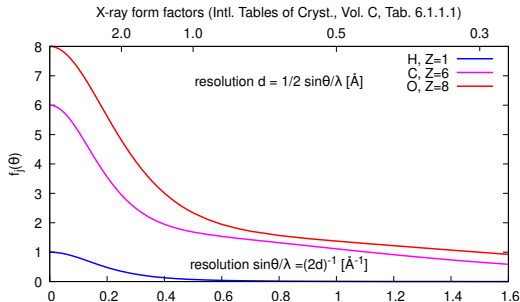
$f_j$  atomic *form* factor. Dependent on atom element, decreases with decreasing scattering angle  $\theta = \theta(hkl)$

$U_j(\theta, \lambda)$  atomic displacement parameter (ADP, alias Debye-Waller factor): models thermal vibration of atoms

$e^{2\pi i(hx_j + ky_j + lz_j)}$  phase shift of the atom relative to the origin of the unit cell

## The form factor $f_j(\theta)$

Intensity of the scattered X-rays decreases with increasing scattering angle  $\theta$ .



The wavelength  $\lambda$  is of the same order of magnitude as the size of the atoms: photons “see” the shape of atoms.

Note: hydrogen atoms do not contribute to data higher than 1 Å

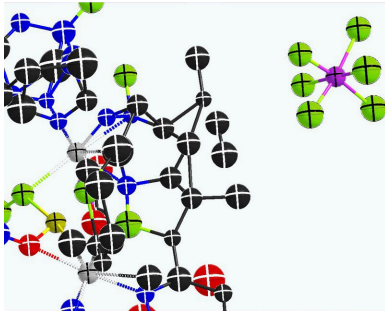
## The atomic displacement parameter $U_j(\theta, \lambda)$

- Atoms vibrate at  $T > 0K$
- Vibration leads to reduction of spot intensities, but not to change in spot shape
- At medium resolution: 1 parameter
- At high resolution: anisotropic description with 6 parameters as ellipsoids

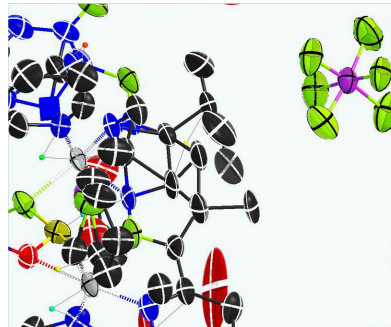
<b>Isotropic</b>		<b>Anisotropic</b>
$U_j(\theta, \lambda) = 4U_{\text{iso}} \frac{\sin^2 \theta}{\lambda^2}$	$U_j(\theta, \lambda) = (hkl) \begin{pmatrix} U_{11} & U_{12} & U_{13} \\ U_{12} & U_{22} & U_{23} \\ U_{13} & U_{23} & U_{33} \end{pmatrix} \begin{pmatrix} a^* \\ b^* \\ c^* \end{pmatrix}$	$\begin{pmatrix} a^* \\ b^* \\ c^* \end{pmatrix}$
1 parameter per atom		6 parameters per atom

- Name : ADP = isotropic or anisotropic *atomic displacement parameter*

## Example images for ADP



Refinement with isotropic ADPs

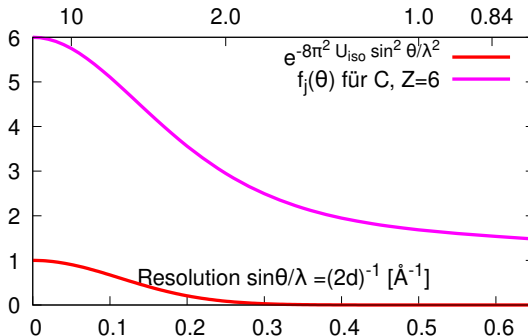


Refinement with anisotropic ADPs

Caveat: Anisotropic ADPs “wipe up” errors in the data.



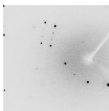
## The ADP $U$ : a fudge factor



- Sharp drop-off with resolution: can make wrongly placed atoms disappear.
- Similarity with form factor: confusion of atom types
- Risk of overfitting

# 5 Phasing

Data collection



several GB



Data integration

0	0	-1	2.7	0.9
0	0	1	4.0	1.0
0	0	-2	1'257.0	35.5
0	0	-2	1'600.0	42.7

several files, 100's MB



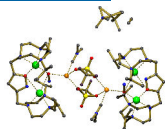
Data Scaling

0	0	-1	2.8	0.55
0	0	1	3.8	0.63
0	0	-2	1'432.0	95.7
0	0	-2	1'282.0	85.9

1 "hkl"-file, 50MB



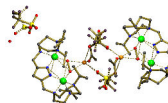
Phasing



Starting model



Refinement



Chemically  
sensible  
model

## Phasing *alias* Solving the structure

- Ideally, we would measure data, run a single calculation, and get the molecular structure as result.
- our data: thousands of measured intensities  $I(hkl)$ .
- From Eq. 2 (p. 40) and Eq. 4 (page 45):

$$\begin{aligned} I(hkl) &= c |F(hkl)|^2 \\ &= c \left| \int_{\text{unit cell}} \rho(x, y, z) e^{2\pi i(hx+ky+lz)} d^3x \right|^2 \end{aligned}$$

- **If** we could invert this equation, we could calculate the coordinates  $x, y, z$  in one go. But we cannot ...

## The phase problem

The inverse of the Fourier transformation

$$F(hkl) = \int_{\text{unit cell}} \rho(x, y, z) e^{2\pi i(hx+ky+lz)} d^3x$$

reads

$$\rho(x, y, z) = \sum_{(h,k,l)} F(hkl) e^{-2\pi i(hx+ky+lz)}$$

## The phase problem

The structure factor  $F(hkl)$  is a complex number. Therefore, it has

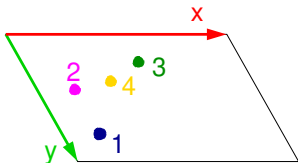
**an amplitude**  $|F(hkl)| = \sqrt{I(hkl)/c}$

**a phase**  $\phi(hkl) = ???$

$$F(hkl) = \sqrt{I(hkl)/c} \times e^{-i\phi(hkl)}$$

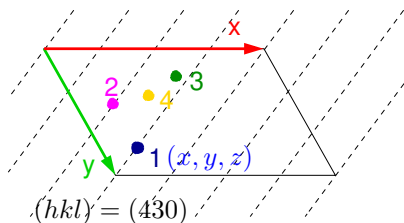
We can measure the amplitude, but we cannot measure the phase. This is known as the **phase problem of crystallography**.

## Relation between Phases and Molecule



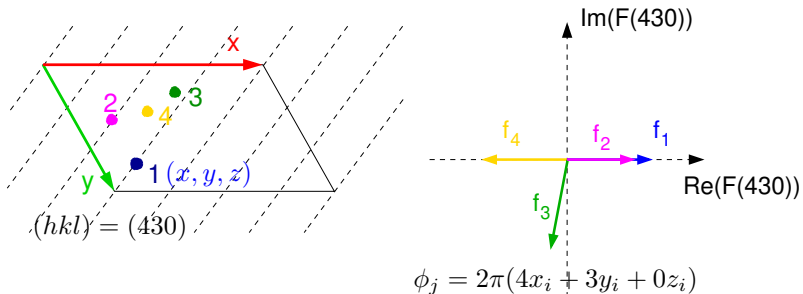
Unit cell with four atoms inside.

## The phases are related to the chemical structure



Example reflection  $F(4\bar{3}0) = \sum_{j=1}^4 f_j(\theta) e^{-8\pi^2 U_j(\theta, \lambda)} e^{2\pi i(4x_j + 3y_j + 0z_j)}$

## The phases are related to the chemical structure



$$\phi_j = 2\pi(4x_i + 3y_i + 0z_i)$$

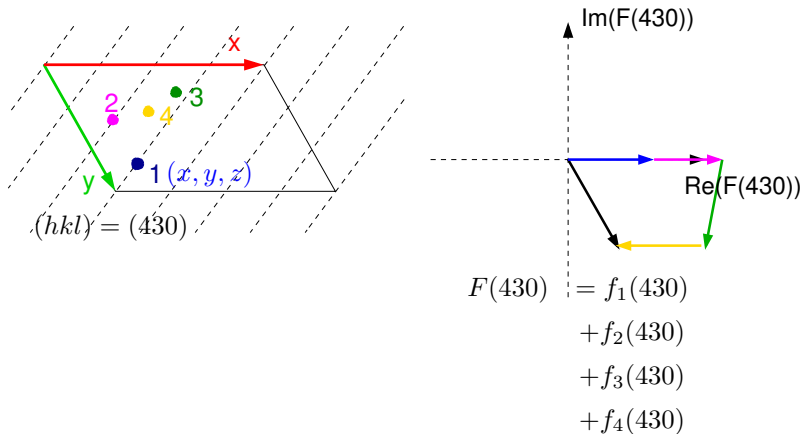
= rel. distance to origin

= rel. distance to lattice plane

$$F(430) = \sum_{j=1}^4 f_j(\theta) e^{-8\pi^2 U_j(\theta, \lambda)} e^{2\pi i(4x_j + 3y_j + 0z_j)}$$

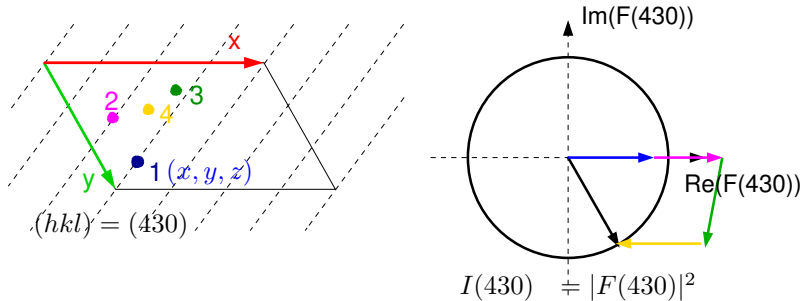


## The phases are related to the chemical structure



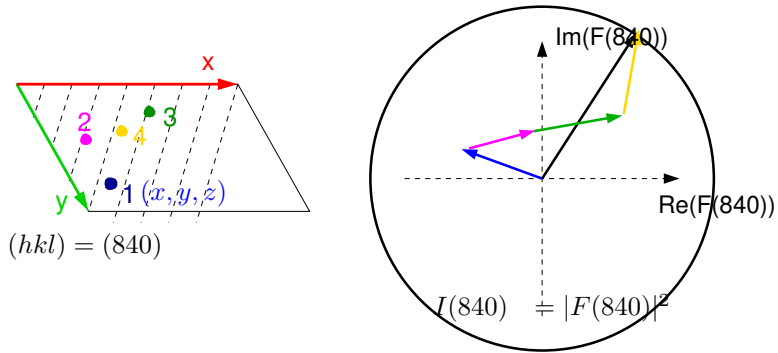
Total  $F(430)$  = sum of individual vectors

## The phases are related to the chemical structure



Measurement  $I(430) = |F(430)|^2$ : loss of phase information

## The phases are related to the chemical structure



Same coordinates, *different* contributions per atom to  $F(840)$

## Phases $\phi(hkl)$

- The phase problem arises, because  $F(hkl)$  is a complex number.
- $F(hkl) = \sum_{\text{all atoms}} f_j(\theta) e^{-8\pi^2 U_j(\theta, \lambda)} e^{2\pi i(hx_j + ky_j + lz_j)}$
- The length  $f_j$  for each atom is independent of  $(hkl)$ .
- The phase contribution  $e^{2\pi i(hx_j + ky_j + lz_j)}$  varies for each reflection.
- the phases  $\phi(hkl)$  is an angle between  $0^\circ$  and  $360^\circ$
- The total phase  $\phi(hkl)$  contains convoluted information from each atom. Each contribution depends on the position of the atom **relative** to the lattice planes for the reflection  $(hkl)$ , *i.e.*, there are as many phases as there are reflection spots.

## References

- [1] Zbigniew Dauter. 'Collection of X-Ray Diffraction Data from Macromolecular Crystals'. In: *Protein Crystallography: Methods and Protocols*. Ed. by Alexander Wlodawer, Zbigniew Dauter and Mariusz Jaskolski. New York, NY: Springer New York, 2017, pp. 165–184. DOI: 10.1007/978-1-4939-7000-1\_7.
- [2] C. Giacovazzo, ed. *Fundamentals of Crystallography*. Oxford University Press, 1985.