

### Chemical Crystallography and Structural Chemistry

VO 270287

### Lecture N° 3 — $18^{th}$ March 2021

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## Announcement: Fridays for Future Austria



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# Reminder: Start the recording!

Tim G	rüne Chemical Crystallography	() universität Wien
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# 1 Diffraction Theory — Atoms and X-rays



### Physicists' description of light

planar wave

spherical wave



$$A(\vec{x},t) = A_0 \sin(\vec{k}\vec{x} - \omega t) \qquad \qquad A(\vec{x},t) = \frac{A_0}{|\vec{x}|}\sin(|k||\vec{x}| - \omega t)$$

- Intensity  $I \propto {A_0}^2$  and  $I \propto (A_0/|x|)^2$ , respectively
- Oscillation  $(\omega t)$  is not observed on the detector, intensity I constant does not "flicker" on the detector
- Direction of propagation:  $\vec{k};~|\vec{k}|=2\pi/\lambda=\omega/c$

Important features of waves: the sum of two waves result in a wave



### Crystal as amplifier



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### Crystal as amplifier





### Crystal as amplifier





### Path difference $\Delta$ (red - blue)

- Each point on the detector results in a specific path difference
- The signal at such point depends on the path difference
- The path difference can be
  - 1. an arbitrary multiple of the wavelength  $\boldsymbol{\lambda}$
  - 2. an integer multiple of the wavelength  $\boldsymbol{\lambda}$
  - 3. an integer multiple + one half of the wavelength  $\lambda$



### Path difference $\Delta$ (red - blue) — arbitrary path difference

Summed wave -----

10 waves, random path difference



Total amplitude is a of same order as individual waves. With many atoms: signal buried in the noise



### Path difference $\Delta$ (red - blue) — multiple of wavelength $\lambda$

Summed wave -----

10 waves, path difference  $n^*\lambda$ 



Total amplitude = 10 \* individual amplitude: regular order amplifies signal



### Path difference $\Delta$ (red - blue) = $(n + 1/2) * \lambda$

Summed wave -----

10 waves, path difference  $(n+1/2)^*\lambda$ 



In certain circumstances there can be complete extinction of the signal. This is important for space group determination.

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#### Crystal as wave amplifier



With constructive interference (right), the amplitude of the signal grows with the number of waves (unit cells), much faster than for random interference (left).



### Regularity of the crystal



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### **Diffraction spots**



- Spot on the detector surface = constructive interference from all unit cells of the crystal
- noise: everywhere else
- Path difference \(\Delta\) (red blue) depends on:
  - 1. direction of incoming ray
  - 2. direction of outgoing ray
  - 3. wavelength
  - periodicity of the crystal = unit cell parameters



### 2 The Laue equations



- A reflection occurs at the detector, where the path difference is an integer multiple of the wave length.
- The locations lie on rays coming from the crystal
- The directions depend on unit cell parameters and crystal orientation
- The directions are described by the Laue equations (Max von Laue, 1879 1960) .



$$\begin{split} \vec{a} \cdot \vec{S} &= |\vec{a}| |\vec{S}| \cos(\vec{a}, \vec{S}) = h \\ \vec{b} \cdot \vec{S} &= |\vec{b}| |\vec{S}| \cos(\vec{b}, \vec{S}) = k \\ \vec{c} \cdot \vec{S} &= |\vec{c}| |\vec{S}| \cos(\vec{c}, \vec{S}) = l \end{split}$$





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The Laue equations describe the geometry of the diffraction experiment:

 $ec{a},ec{b},ec{c}$ : Orientation of the crystal $|ec{S}_{\sf in}|=1/\lambda$ : wavelength of the experiment

 $|\vec{S}_{out}|$ : direction, *alias* position at the detector (is there a spot or not?)

h,k,l integer: integer multiple of path differences  $\Delta=n\cdot\lambda$ 



- Each scattering vector  $\vec{S}$  describes exactly one position on the detector
- Only those positions, that fulfil all three Laue equations at once, will show a reflection
- Each reflection is uniquely described by the triplet of integers  $\left(hkl\right)$
- The triplet (h, k, l) is called **the Miller index** of the corresponding reflection (W. H. Miller, 1801–1880)
- The direct beam  $\vec{S}_{\rm in}$  coincides with the reflection  $(0,0,0),~\vec{S}_{\rm in}=\vec{S}_{\rm out}$
- The reflection  $\left(0,0,0\right)$  cannot be measured!



### Indexing

- the term **indexing** describes the assignment of the Miller indices to and the reflections recorded on the detector.
- Indexing is equivalent to determining the unit cell parameters a,b,c,  $\alpha,\beta,\gamma$  and the crystal orientation.
- Indexing is an essential step for data processing



$$\leftrightarrow$$

$$\begin{split} \vec{a} \cdot \vec{S} &= |\vec{a}| |\vec{S}| \cos(\vec{a}, \vec{S}) = h \\ \vec{b} \cdot \vec{S} &= |\vec{b}| |\vec{S}| \cos(\vec{b}, \vec{S}) = k \\ \vec{c} \cdot \vec{S} &= |\vec{c}| |\vec{S}| \cos(\vec{c}, \vec{S}) = l \end{split}$$



### Indexing

In most cases, indexing algorithms work very easily and fast Reasons for difficulties with indexing:

- More than one crystal lattice (twinning)
- Very large unit cell leads to overlapping reflections

 $\leftrightarrow$ 

• Wrong values for the experimental parameters (detector distance, wavelength, rotation axis)



$$\begin{split} \vec{a} \cdot \vec{S} &= |\vec{a}| |\vec{S}| \cos(\vec{a}, \vec{S}) = h \\ \vec{b} \cdot \vec{S} &= |\vec{b}| |\vec{S}| \cos(\vec{b}, \vec{S}) = k \\ \vec{c} \cdot \vec{S} &= |\vec{c}| |\vec{S}| \cos(\vec{c}, \vec{S}) = l \end{split}$$

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### The Ewald Sphere

Laue equations: mathematically handy, but difficult to imagine

**Ewald sphere:** Construction to help understand the diffraction pattern

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### Prelude: The reciprocal lattice and the Miller Indices

- The corners of the unit cell span the crystal lattice
- The unit cell vectors  $\vec{a}, \vec{b}, \vec{c}$  build the basis for the crystal lattice
- Many aspects of diffraction can be described more easily with the help of the **reciprocal lattice**.

• 
$$\vec{a}^* = \frac{\vec{b} \times \vec{c}}{V}$$
:  $\vec{a}^* \perp \text{plane}(\vec{b}, \vec{c})$ 

• 
$$\vec{b}^* = \frac{\vec{c} \times \vec{a}}{V}$$
:  $\vec{b}^* \perp \text{plane}(\vec{c}, \vec{a})$ 

• 
$$\vec{c}^* = \frac{\vec{a} \times \vec{b}}{V}$$
:  $\vec{c}^* \perp \text{plane}(\vec{a}, \vec{b})$ 

V: unit cell volume (Å<sup>3</sup>),  $V = (\vec{a} \times \vec{b}) \cdot c$ 



### Prelude: The reciprocal lattice and the Miller Indices

• 
$$\vec{a}^* = \frac{\vec{b} \times \vec{c}}{V}$$
:  $\vec{a}^* \perp \text{plane}(\vec{b}, \vec{c})$ 

•  $\vec{b}^* = \frac{\vec{c} \times \vec{a}}{V}$ :  $\vec{b}^* \perp \text{plane}(\vec{c}, \vec{a})$  V: unit cell volume (Å<sup>3</sup>),  $V = (\vec{a} \times \vec{b}) \cdot c$ 

•  $\vec{c}^* = \frac{\vec{a} \times \vec{b}}{V}$ :  $\vec{c}^* \perp \text{plane}(\vec{a}, \vec{b})$ The Miller indices (hkl) span the reciprocal lattice:

$$h\vec{a}^* + k\vec{b}^* + l\vec{c}^* \qquad \qquad h, k, l \in \mathbb{Z}$$

- Each reflection corresponds to one point of the reciprocal lattice.
- When  $\alpha = \beta = \gamma = 90^{\circ}$  (orthorhombic unit cell):  $|\vec{a}^*| = 1/a$ ,  $|\vec{b}^*| = 1/b$ ,  $|\vec{c}^*| = 1/c$















Draw a sphere with radius  $1/\lambda$  that **touches** the lattice origin. The sphere centre lies aligned with the X-ray source.

This sphere is the **Ewald Sphere**.





The scattering vector  $\vec{S}$  points from the origin to the lattice point.

Some lattice points **touch** the surface of the Ewald sphere (red circles). These fulfil the Laue conditions.

They are the **recordable** reflections.









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## 4 Bragg's Law



Idea:

- X-rays are reflected on lattice planes (German: "Gitterebenen" or "Net-zebenen")
- Reflections occur when the path difference is a multiple integer of the wave length



1. Lattice: Corners of the unit cells.



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2. Three corner points make a plane.





1. Lattice: Corners of

### Bragg's Law







- 1. Lattice: Corners of the unit cells.
- 2. Three corner points make a plane.
- 3. Parallel shifting of plane through **all** lattice corners creates a set of planes.
- 4. Constructive interferences leads to Bragg's Law:

 $n\lambda=2d\sin\theta$ 

d : distance between planes.

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### Bragg's Law and Resolution of a Reflection

 $n\lambda=2d\sin\theta$ 

- The value *d* is called **the resolution** of the reflection (*hkl*)
- *d* is measured in Å.
- N.B.: High resolution corresponds to a small value of  $\boldsymbol{d}$
- $n \ge 2$  higher order reflections, which usually do not occur. We only need to consider the case n = 1, *i.e.*

 $\lambda=2d\sin\theta$ 



### Bragg's Law and Resolution of a Reflection

 $\lambda=2d\sin\theta$ 

Sometimes, the inverted value

$$d^* \equiv 1/d = \frac{2\sin\theta}{\lambda}$$

is called resolution, measured in  $1/\textrm{\AA}$ 

- High values of  $d^*$  correspond to high resolution and vice versa
- $d^*$  has the length of the reciprocal lattice vector

$$d^* = 1/d = \|h\vec{a}^* + k\vec{b}^* + l\vec{c}^*\|$$



### Example image with resolution rings



- All reflections on a circle about the direct beam have the same resolution
- When the detector if offset  $(2\theta \neq 0^{\circ})$ , the circles become ellipses (intersection of a cone and a plane)



# **Reflection Intensity**

- The Laue equations, the Bragg equation, and the Ewald sphere construction all refer to the spot positions.
- They contain no information about spot intensity
- The molecule inside the unit cell determines the intensity of every reflections:

$$\begin{split} I(hkl) \propto |\sum_{\text{atom } j} f_j \cos(hx_j + ky_j + lz_j)|^2 \\ \propto |\sum_{\text{atom } j} f_j e^{2\pi i (hx_j + ky_j + lz_j)}|^2 \end{split}$$

- 1. every atom emits a small spherical wave (cos-term)
- 2. the amplitude  $f_j$  is proportional to the atom number Z
- 3.  $f_j$  is called *atomic scattering factor*



### Summary Laue equations

- Reflex positions ("patterns") depend on the unit cell parameters  $a,b,c,\alpha,\beta,\gamma$  and the orientiation of the crystal.
- Reflex positions *do not* depend on the chemical content of the unit cell.
- Reflex intensities depend on the chemical content of the unit cell
- Ewald sphere visualises the Laue equation
- Every reflex has a resolution *d*, via Bragg's law.
- High resolution = small d, low resolution = large d